

Langevin Dynamics from Microscopic Model [nln74]

Brownian particle harmonically coupled to N otherwise free particles that serve as a primitive form of heat bath. [Wilde and Singh 1998]

Classical Hamiltonian:

$$\mathcal{H} = \frac{p^2}{2m} + \sum_{i=1}^N \left(\frac{p_i^2}{2m_i} + \frac{1}{2} m_i \omega_i^2 \bar{x}_i^2 \right), \quad \bar{x}_i \doteq x_i - \frac{c_i x}{m_i \omega_i^2}, \quad (1)$$

where $m_i \omega_i^2$ is the stiffness of the harmonic coupling between the Brownian particle and one of the heat-bath particles. The c_i are conveniently scaled coupling constants.

Canonical equations:

$$\frac{dx}{dt} = \frac{\partial \mathcal{H}}{\partial p} = \frac{p}{m}, \quad \frac{dp}{dt} = -\frac{\partial \mathcal{H}}{\partial x} = \sum_{i=1}^N c_i \bar{x}_i, \quad (2a)$$

$$\frac{dx_i}{dt} = \frac{\partial \mathcal{H}}{\partial p_i} = \frac{p_i}{m_i}, \quad \frac{dp_i}{dt} = -\frac{\partial \mathcal{H}}{\partial x_i} = -m_i \omega_i^2 \bar{x}_i; \quad i = 1, \dots, N. \quad (2b)$$

Elimination of momenta yields 2nd-order ODEs:

$$m \frac{d^2 x}{dt^2} = m \frac{d \dot{x}}{dt} = F(t), \quad F(t) \doteq \sum_{i=1}^N c_i \bar{x}_i, \quad (3a)$$

$$m_i \frac{d^2 x_i}{dt^2} = -m_i \omega_i^2 x_i + c_i x, \quad i = 1, \dots, N. \quad (3b)$$

Formal solution of (3b):

$$x_i(t) = x_i(0) \cos(\omega_i t) + \frac{\dot{x}_i(0)}{\omega_i} \sin(\omega_i t) + \underbrace{\frac{c_i}{m_i \omega_i} \int_0^t dt' x(t') \sin(\omega_i(t-t'))}_{A(t)}. \quad (4)$$

Integrate by parts:

$$A(t) = \frac{1}{\omega_i} \left[x(t) - x(0) \cos(\omega_i t) - \int_0^t dt' \dot{x}(t') \cos(\omega_i(t-t')) \right]. \quad (5)$$

Assemble parts, then use (1) and (3a):

$$\begin{aligned} x_i(t) &= \left(x_i(0) - \frac{c_i}{m_i \omega_i^2} x(0) \right) \cos(\omega_i t) + \frac{\dot{x}_i(0)}{\omega_i} \sin(\omega_i t) \\ &\quad + \frac{c_i}{m_i \omega_i^2} \left[x(t) - \int_0^t dt' \dot{x}(t') \cos(\omega_i(t-t')) \right], \end{aligned} \quad (6)$$

$$\bar{x}_i(t) = \bar{x}_i(0) \cos(\omega_i t) + \frac{\dot{x}_i(0)}{\omega_i} \sin(\omega_i t) - \frac{c_i}{m_i \omega_i^2} \int_0^t dt' \dot{x}(t') \cos(\omega_i(t-t')),$$

$$F(t) = \sum_{i=1}^N \left[c_i \bar{x}_i(0) \cos(\omega_i t) + \frac{c_i}{\omega_i} \dot{x}_i(0) \sin(\omega_i t) - \frac{c_i^2}{m_i \omega_i^2} \int_0^t dt' \dot{x}(t') \cos(\omega_i(t-t')) \right]. \quad (7)$$

Expectation values at thermal equilibrium:

$$\langle \bar{x}_i(0) \rangle = 0, \quad \langle \dot{x}_i(0) \rangle = 0, \quad \omega_i^2 \langle \bar{x}_i(0) \bar{x}_j(0) \rangle = \langle \dot{x}_i(0) \dot{x}_j(0) \rangle = \frac{k_B T}{m_i} \delta_{ij}.$$

$$\langle F(t) \rangle = - \int_0^t dt' \dot{x}(t') \alpha_s(t-t'), \quad \alpha_s(t-t') = \underbrace{\sum_{i=1}^N \frac{c_i^2}{m_i \omega_i^2} \cos(\omega_i(t-t'))}_{\text{attenuation function}}.$$

Random force:

$$f(t) \doteq F(t) - \langle F(t) \rangle = \sum_{i=1}^N \left[c_i \bar{x}_i(0) \cos(\omega_i t) + \frac{c_i}{\omega_i} \dot{x}_i(0) \sin(\omega_i t) \right]. \quad (8)$$

Generalized Langevin equation:

$$m \frac{d\dot{x}}{dt} = - \int_0^t dt' \dot{x}(t') \alpha_s(t-t') + f(t). \quad (9)$$

Fluctuation-dissipation relation:

$$\begin{aligned} \langle f(t) f(t') \rangle &= \sum_{i=1}^N \left[c_i^2 \cos(\omega_i t) \cos(\omega_i t') \underbrace{\langle \bar{x}_i(0) \bar{x}_i(0) \rangle}_{k_B T / m_i \omega_i^2} \right. \\ &\quad \left. + \frac{c_i^2}{\omega_i^2} \sin(\omega_i t) \sin(\omega_i t') \underbrace{\langle \dot{x}_i(0) \dot{x}_i(0) \rangle}_{k_B T / m_i} \right] \\ &= k_B T \sum_{i=1}^N \frac{c_i^2}{m_i \omega_i^2} \cos(\omega_i(t-t')) = k_B T \alpha_s(t-t'). \end{aligned} \quad (10)$$