## Bessel Functions [gmd4F]

Solutions of partial differential equations for problems with cylindrical symmetry are often expressible as Bessel functions.

Bessel equation: 
$$z^2 R''(z) + z R'(z) + (z^2 - \nu^2) R(z) = 0.$$
  
Bessel functions of the first kind:  $J_{\nu}(z) = \sum_{s=0}^{\infty} \frac{(-1)^s}{s! \Gamma(s+\nu+1)} \left(\frac{z}{2}\right)^{2s+\nu}$ 

For noninteger order  $\nu$ , the function  $J_{\nu}(z)$  and  $J_{-\nu}(z)$  are linearly independent. Their linear dependence for integer  $\nu$  is manifest in the relation,<sup>1</sup>



$$J_{-\nu}(z) = (-1)^{\nu} J_{\nu}(z) \quad : \ \nu \in \mathbb{Z}.$$

A solution which remains linearly independent of  $J_{\nu}(x)$  is the Neumann function (Bessel function of the second kind) constructed as follows:



The divergence at x = 0 is logarithmic in nature.

<sup>&</sup>lt;sup>1</sup>This relation emerges on account of the fact that the function  $\Gamma(x)$  diverges for nonpositive integers.

## Modified Bessel functions:

A  $90^{\circ}$  rotation in the complex plane converts the Bessel functions into modified versions such as the followig:

$$\mathbf{I}_{\nu}(z) = \imath^{-\nu} \mathbf{J}_{\nu}(\imath z).$$

They are solutions of the modified Bessel equation,

$$z^{2}R''(z) + zR'(z) - (z^{2} + \nu^{2})R(z) = 0,$$

and have quite different properties for real z.

Modified Bessel function of the first kind:  $I_{\nu}(z) = \sum_{s=0}^{\infty} \frac{1}{s!\Gamma(s+\nu+1)} \left(\frac{z}{2}\right)^{2s+\nu}$ . For noninteger  $\nu$ , the functions  $I_{\nu}(z)$  and  $I_{-\nu}(z)$  are again linearly indepen-

dent, whereas for integer  $\nu$  we have,

$$\mathbf{I}_{-\nu}(z) = \mathbf{I}_{\nu}(z) \quad : \ \nu \in \mathbb{Z}.$$

A solution which remains linearly independent of  $I_{\nu}(z)$  for integer  $\nu$  is the *MacDonald function* (modified Bessel function of the second kind):

$$\mathbf{K}_{\nu}(z) \doteq \frac{\pi}{2} \frac{\mathbf{I}_{-\nu}(z) - \mathbf{I}_{\nu}(z)}{\sin(\nu\pi)}.$$



Useful relations:

$$\begin{split} & \triangleright \ \frac{d}{dz} \left[ z^{\nu} \mathbf{J}_{\nu}(z) \right] = z^{\nu} \mathbf{J}_{\nu-1}(z), \quad \frac{d}{dz} \left[ z^{-\nu} \mathbf{J}_{\nu}(z) \right] = -z^{-\nu} \mathbf{J}_{\nu+1}(z). \\ & \triangleright \ \mathbf{J}_{\nu}'(z) = \mathbf{J}_{\nu-1}(z) - \frac{\nu}{z} \mathbf{J}_{\nu}(z), \quad \mathbf{J}_{\nu}'(z) = -\mathbf{J}_{\nu+1}(z) + \frac{\nu}{z} \mathbf{J}_{\nu}(z). \\ & \triangleright \ \mathbf{J}_{\nu}'(z) = \frac{1}{2} \left[ \mathbf{J}_{\nu-1}(z) - \mathbf{J}_{\nu+1}(z) \right]. \\ & \triangleright \ \mathbf{J}_{\nu}(z) = \frac{1}{2} \nu z \left[ \mathbf{J}_{\nu-1}(z) + \mathbf{J}_{\nu+1}(z) \right]. \\ & \triangleright \ \mathbf{J}_{0}(z) = \frac{1}{2\pi} \int_{0}^{2\pi} d\theta \, e^{iz\cos\theta}. \\ & \triangleright \ |z| \ll 1: \quad \mathbf{J}_{\nu}(z) \rightsquigarrow \frac{(z/2)^{\nu}}{\Gamma(\nu+1)} \quad : \nu \neq 0, -1, -2, \dots, \\ & \text{power-law rise from zero or power-law divergence.} \\ & \triangleright \ |z| \gg 1: \quad \mathbf{J}_{\nu}(z) \rightsquigarrow \sqrt{\frac{2}{\pi z}} \cos\left(z + \left(\nu + \frac{1}{2}\right)\frac{\pi}{2}\right), \\ & \text{attenuated oscillation with $\nu$-dependent phase shift.} \end{split}$$