## Bessel Functions [gmd4F]

Solutions of partial differential equations for problems with cylindrical symmetry are often expressible as Bessel functions.

Bessel equation: 
$$
z^2 R''(z) + z R'(z) + (z^2 - \nu^2) R(z) = 0
$$
.  
\nBessel functions of the first kind:  $J_{\nu}(z) = \sum_{s=0}^{\infty} \frac{(-1)^s}{s! \Gamma(s + \nu + 1)} \left(\frac{z}{2}\right)^{2s + \nu}$ 

For noninteger order  $\nu$ , the function  $J_{\nu}(z)$  and  $J_{-\nu}(z)$  are linearly independent. Their linear dependence for integer  $\nu$  is manifest in the relation,<sup>1</sup>

.



$$
J_{-\nu}(z) = (-1)^{\nu} J_{\nu}(z) \quad : \nu \in \mathbb{Z}.
$$

A solution which remains linearly independent of  $J_{\nu}(x)$  is the *Neumann func*tion (Bessel function of the second kind) constructed as follows:



The divergence at  $x = 0$  is logarithmic in nature.

<sup>&</sup>lt;sup>1</sup>This relation emerges on account of the fact that the function  $\Gamma(x)$  diverges for nonpositive integers.

## Modified Bessel functions:

A 90 $^{\circ}$  rotation in the complex plane converts the Bessel functions into modified versions such as the folowig:

$$
I_{\nu}(z) = i^{-\nu} J_{\nu}(iz).
$$

They are solutions of the modified Bessel equation,

$$
z2R''(z) + zR'(z) - (z2 + \nu2)R(z) = 0,
$$

and have quite different properties for real z.

Modified Bessel function of the first kind:  $I_{\nu}(z) = \sum_{n=0}^{\infty}$  $s=0$ 1  $s!\Gamma(s+\nu+1)$  $\frac{z}{2}$ 2  $\big)^{2s+\nu}$ . For noninteger  $\nu$ , the functions  $I_{\nu}(z)$  and  $I_{-\nu}(z)$  are again linearly indepen-

dent, whereas for integer  $\nu$  we have,

$$
I_{-\nu}(z) = I_{\nu}(z) \quad : \nu \in \mathbb{Z}.
$$

A solution which remains linearly independent of  $I_{\nu}(z)$  for integer  $\nu$  is the MacDonald function (modified Bessel function of the second kind):

$$
K_{\nu}(z) \doteq \frac{\pi}{2} \frac{I_{-\nu}(z) - I_{\nu}(z)}{\sin(\nu \pi)}.
$$



## Useful relations:

$$
\triangleright \frac{d}{dz} [z^{\nu} J_{\nu}(z)] = z^{\nu} J_{\nu-1}(z), \quad \frac{d}{dz} [z^{-\nu} J_{\nu}(z)] = -z^{-\nu} J_{\nu+1}(z).
$$
  
\n
$$
\triangleright J_{\nu}'(z) = J_{\nu-1}(z) - \frac{\nu}{z} J_{\nu}(z), \quad J_{\nu}'(z) = -J_{\nu+1}(z) + \frac{\nu}{z} J_{\nu}(z).
$$
  
\n
$$
\triangleright J_{\nu}'(z) = \frac{1}{2} [J_{\nu-1}(z) - J_{\nu+1}(z)].
$$
  
\n
$$
\triangleright J_{\nu}(z) = \frac{1}{2} \nu z [J_{\nu-1}(z) + J_{\nu+1}(z)].
$$
  
\n
$$
\triangleright J_{0}(z) = \frac{1}{2\pi} \int_{0}^{2\pi} d\theta e^{iz \cos \theta}.
$$
  
\n
$$
\triangleright |z| \ll 1: \quad J_{\nu}(z) \rightsquigarrow \frac{(z/2)^{\nu}}{\Gamma(\nu+1)} \quad : \quad \nu \neq 0, -1, -2, \dots,
$$
  
\npower-law rise from zero or power-law divergence.  
\n
$$
\triangleright |z| \gg 1: \quad J_{\nu}(z) \rightsquigarrow \sqrt{\frac{2}{\pi z}} \cos\left(z + (\nu + \frac{1}{2})\frac{\pi}{2}\right),
$$
  
\nattenuated oscillation with  $\nu$ -dependent phase shift.